

# Searches for Neutral Higgs Boson and Interpretations in the MSSM at LEP

Philip Bechtle<sup>a</sup>

DESY, Notkestr. 85, 22607 Hamburg, Germany; E-mail: philip.bechtle@desy.de

Received: 1 Oct 2003 / Accepted: 20 Nov 2003 /

Published Online: 26 Nov 2003 – © Springer-Verlag / Società Italiana di Fisica 2003

**Abstract** This paper discusses recent publications of the LEP collaborations DELPHI, L3 and OPAL on searches for Higgs bosons motivated by MSSM scenarios as well as their interpretation in the MSSM. With the final publication of the LEP collaborations available or awaited, more and more interpretations in different MSSM models, including both CP conserving and CP violating, become available. Also specialized analyses close open areas in the parameter space. In the same time, better theoretical calculations with an increased maximal mass of the  $h$  boson were presented. Both the new scenarios as well as the new theoretical limit on  $m_h$  has consequences for the limits from LEP. The searches, the models in which they are interpreted and the implications of the LEP results for future SUSY searches, especially on the  $\tan \beta$  limit, are presented here.

**PACS.** 12.60.Fr Extensions of electroweak Higgs sector – 12.60.Jv Supersymmetric models – 13.66.Fg Gauge and Higgs boson production in  $e^+e^-$  interactions – 14.80.Cp Non-standard-model Higgs bosons

## 1 Introduction

In the Standard Model (SM) it is generally assumed that the Higgs mechanism is responsible for the breaking of electroweak symmetry and for the generation of elementary particle masses. The Minimal Supersymmetric Standard Model (MSSM) is the SUSY extension of the SM with minimal new particle content. It introduces two complex Higgs field doublets. The MSSM predicts five Higgs bosons: three neutral and two charged ones. At least one of the neutral Higgs bosons is predicted to have its mass close to the electroweak energy scale, providing a high motivation to the searches at current and future colliders.

In the MSSM the Higgs potential is assumed to be invariant under CP transformation at tree level. However, it is possible to break CP symmetry in the Higgs sector by radiative corrections, especially by contributions from complex trilinear couplings  $A_{t,b}$  of third generation scalar-quarks [1].

Since the input parameter space is generally too large to be scanned completely, so called benchmark scenarios (cf. Tab. 1) have been proposed [2, 3, 4], each emphasising a certain phenomenological situation. The parameters  $\tan \beta = v_2/v_1$  and  $m_A$  governing the Higgs sector on tree-level are scanned, while all parameters on loop level are kept fixed for one scenario. CP conserving (CPC) and CP violating (CPV) scenarios exist.

Depending on the parameters of the MSSM, Higgs Bosons can be produced in Higgstrahlung  $e^+e^- \rightarrow hZ$ , as in the SM, or in pair production  $e^+e^- \rightarrow hA$ . Flavour independent Higgs decays, Higgs decays into invisible particles or decays of the type  $h \rightarrow AA$  additionally present new topologies.

**Table 1.** The MSSM scenarios used by the LEP collaborations, proposed in [2], [3] and [4]

CP conserving	
No Mixing	No mixing in the stop-sbottom sector
$m_h$ max	Maximum $m_h$ for given $\tan \beta$ , $m_A$
Large $\mu$	Always kinematically accessible, but $h \rightarrow b\bar{b}$ suppressed
$m_h$ max <sup>+</sup>	like $m_h$ max, but favoured by $(g-2)_\mu$ .
constr. $m_h$ max	like $m_h$ max, but favoured by $(b \rightarrow s\gamma)$
gluophobic	hgg coupling suppressed, reduced LHC production cross section
small $\alpha_{\text{eff}}$	$h \rightarrow b\bar{b}$ suppressed by cancellation of $\tilde{b} - \tilde{g}$ loops
CP violating	
CPX	Mixing of CP- and mass-eigenstates several derivatives under study

Since the end of LEP running, the LEP collaborations have developed new searches closing some of these unexcluded areas in the parameter space, thereby giving important new information about the tasks left over for future experiments. This publication will focus on new searches dedicated to formerly uncovered final states described in Section 2, on the consequences of new theoretical developments in Section 3 and on the interpretation of the MSSM Higgs searches in the benchmark scenarios in Section 4.

<sup>a</sup> Speaker; on behalf of the LEP collaborations

## 2 Searches for Higgs bosons in the MSSM

The search for Higgs bosons in the MSSM uses a large variety of channels. The SM production channels Higgsstrahlung  $e^+e^- \rightarrow hZ$  and Boson fusion are reinterpreted in the MSSM. Additionally, dedicated searches for Higgsstrahlung channels with Higgs decays in the MSSM exist. Also, pair production  $e^+e^- \rightarrow hA$  and Yukawa production  $e^+e^- \rightarrow b\bar{b}h/A$  channels are used. New searches comprise the search for  $e^+e^- \rightarrow hZ$  with  $h \rightarrow AA$ , with  $m_A < 10$  GeV below the  $b\bar{b}$  production threshold [5] by OPAL. The same final state with heavier  $m_A$  has been sought-after for the interpretation in CPV models [6].

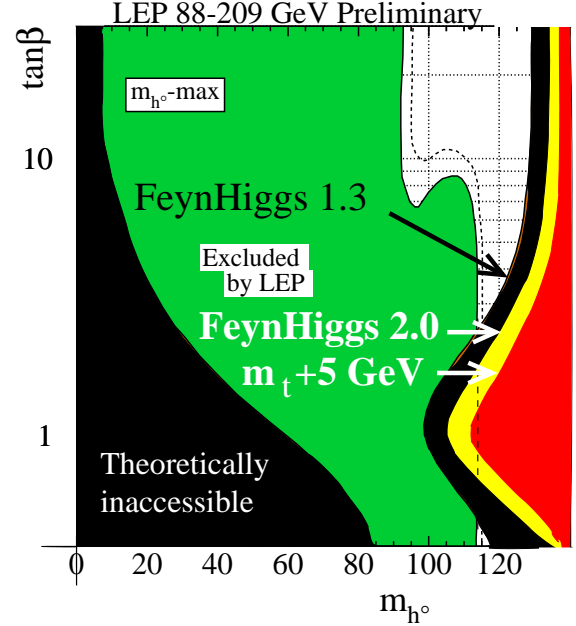
The search from DELPHI [7] for invisibly decaying Higgs bosons has been interpreted in a modified  $m_h$ -max scenario with  $M_2 = \mu = 150$  GeV. It shows that the benchmark scenarios from Tab. 1 do not cover the full range of possible MSSM topologies in the Higgs sector, since the search for invisibly decaying Higgs bosons is needed to cover areas unexcluded by the standard searches.

## 3 New benchmark scenarios and Higgs mass calculations in the MSSM

The calculations of the observables of the Higgs sector, the masses, branching ratios and cross-sections, depending on the choice of SUSY parameters have been performed using two different calculation tools. FEYNHIGGS [8,9] is based on the two-loop diagrammatic approach of [10], and SUBH-POLE/CPH [4] is based on the one-loop renormalization-group improved calculation of [11, 12].

The first three CPC benchmark scenarios of Tab. 1 have traditionally been considered in the past and have now been extended to scenarios 4 and 5, motivated by limits on the branching ratio of the inclusive decay of a B meson into strange particle states and a photon  $B \rightarrow X_s \gamma$  and muon anomalous magnetic moment  $(g-2)_\mu$  measurements. The last two CPC benchmark scans are aiming to set the stage for future Higgs searches at the LHC. There, some of the dominant search channels would be suppressed, resulting in a reduced search sensitivity. The CPV scenario CPX maximises the mixing of CP- and mass eigenstates and has been tested with several different parameter settings [6].

With respect to the calculations used for the MSSM Higgs LEP combination in 2001 [13], new 2-loop calculations of top loop corrections to the Higgs boson mass have become available [14]. They shift the maximal  $m_h$  achievable in the  $m_h$ -max scenario upwards by up to 5 GeV. The maximal  $m_h$  lies at about 135 GeV. While the expected experimental lower limit on  $m_h$  for low  $\tan\beta$  is not expected to change much with respect to the latest LEP exclusion (cf. Fig. 1), the theoretical upper limit on  $m_h$  shifted from the border of the black area to the border of the light grey (yellow) area. If the top mass would additionally shift upwards from its current central value of  $m_t = 174.3 \pm 5.1$  GeV [15] by only one sigma (which could well be the case, given latest measurements from D0 in the leptonic decay channel [16]), then the upper limit on  $m_h$  would increase again by almost 5 GeV, as indicated by the border of the dark grey (red) area in Fig. 1.



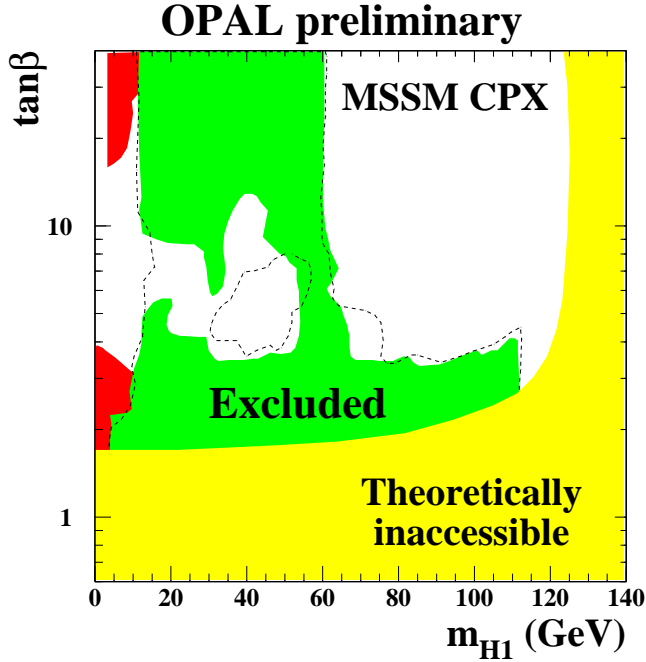
**Figure 1.** Expected exclusion areas of the  $m_h$ -max scenario in the  $m_h$ ,  $\tan\beta$  projection for increased upper bounds on the mass of the h boson. The experimental exclusion areas shown are those of the LEP combination 2001 [13]. The expected theoretical upper bounds for  $m_h$  given new order  $\alpha_t^2$  loop corrections to the Higgs mass is shown in light grey (yellow) and the additional effect of a possible shift in  $m_t$  by  $+1\sigma$  is overlaid in dark grey (red). The theoretically inaccessible area is forbidden by theory.

This example shows that new theoretical developments and a higher precision on  $m_t$  could well influence the exclusion of low  $\tan\beta$  by LEP. A higher precision on  $m_t$  is therefore highly desirable. This could also have implication on the search channels that have to be investigated at future accelerator experiments at LHC searching for a MSSM Higgs boson, where the region of small  $\tan\beta$  can not be regarded as excluded by LEP.

## 4 Interpretation of the Higgs searches in the MSSM

The combination of all Higgs searches of one experiment is used to derive [17,18,6] exclusions in the MSSM parameter space for the scenarios in Tab. 1. No major changes with respect to previous interpretations are received for the no-mixing,  $m_h$ -max and large- $\mu$  scenarios. The latter now can be almost completely excluded by one experiment alone [17], thanks to flavour independent searches [19]. OPAL has also studied [6] the new CPC scenarios 4 to 7 from Tab. 1.

In summary, the parameter choices of the new CPC benchmark scenarios introduce no need for new searches at LEP. Latest in a LEP combination, all topologies are covered up to the kinematic limits of the production channels. The limits on  $m_h$  and  $m_A$  are around 85 to 90 GeV for all CPC scenarios [20, 17, 18,6].



**Figure 2.** Exclusion area of the CPX scenario in the  $\tan\beta, m_{H_1}$  projection from OPAL. The theoretically inaccessible area is forbidden by theory.

The CPX scenario with maximal CP violation in the Higgs sector shows a decoupling of the lightest Higgs bosons  $H_1$  from the Z in the intermediate  $\tan\beta$  range from 4 to 10 ( $H_1$  and  $H_2$  being the lightest and next-to-lightest Higgs boson mass eigenstates). There anyhow  $H_2$  couples to the Z and is heavier than around 100 GeV. Where kinematically accessible, the decay  $H_2 \rightarrow H_1 H_1$  is dominant. Another difference to the common CPC scenarios is the large mass difference  $m_{H_2} - m_{H_1}$  in the range with dominant pair production.

Fig. 2 [6] shows the exclusion areas of the CPX scenario. In the region with dominant  $e^+e^- \rightarrow H_2 Z \rightarrow H_1 H_1 Z$  production at intermediate  $\tan\beta$  and  $m_{H_1} < 50$  GeV open areas emerge. It is expected that a LEP combination will be able to close these holes. Also the lower limit on  $m_{H_1}$  in the large  $\tan\beta$  region, where pair production dominates, is reduced due to the large  $m_{H_2} - m_{H_1}$ . At  $\tan\beta > 5$  and  $m_{H_1} < 10$  GeV, below the  $b\bar{b}$  production threshold and in the pair production region, hardly any experimental constraints exist, since no pair production searches for  $m_{H_2} \approx 100$  GeV and  $m_{H_1} < 10$  GeV exist. Only at large  $\tan\beta > 20$  Yukawa production searches can be used.

## 5 Conclusions

The developments in the MSSM Higgs searches at LEP after the end of LEP data taking in November 2000 exhibit four important lessons. First, also the increased set of CPC and CPV benchmark scenarios do not cover the full range of possible experimental phenomena in the MSSM Higgs sector. Therefore, secondly, a large variety of individual searches is necessary to

cover the rich physics spectrum of the MSSM Higgs sector, which only now become fully available.

Third, new theoretical developments can influence limits on MSSM parameters. Especially the  $\tan\beta$  exclusion of the final LEP combination could be affected. This is also important for the possible MSSM topologies in Higgs searches at the LHC. The importance of external measurements like  $m_t$  from the Tevatron becomes evident. A greater precision on  $m_t$  would be beneficial. Fourth, CPV scenarios show that there is still no strict lower limit on the Higgs mass from LEP. Especially in regions with low  $m_{H_1}$ , but either dominant  $e^+e^- \rightarrow H_2 Z$  or dominant  $e^+e^- \rightarrow H_1 H_2$  production no  $\tan\beta$  independent limit on the Higgs mass exists. Also these regions must probably be sought by future colliders.

## References

1. A. Pilaftsis and C. E. Wagner, Nucl. Phys. B **553** (1999) 3.
2. M. Carena, S. Heinemeyer, C. E. Wagner and G. Weiglein, arXiv:hep-ph/9912223.
3. M. Carena, S. Heinemeyer, C. E. Wagner and G. Weiglein, Eur. Phys. J. C **26** (2003) 601.
4. M. Carena, J. R. Ellis, A. Pilaftsis and C. E. Wagner, Phys. Lett. B **495** (2000) 155.
5. G. Abbiendi *et al.* [OPAL Collaboration], Eur. Phys. J. C **27** (2003) 483.
6. OPAL Collaboration, *Search for Neutral Higgs Bosons Predicted by CP Conserving and CP Violating MSSM Scenarios with the OPAL detector at LEP*, 2003, OPAL PN524
7. [DELPHI Collaboration], *Searches for invisibly decaying Higgs bosons with the DELPHI detector at LEP*, DELPHI 2003-036 CONF 656.
8. S. Heinemeyer, W. Hollik and G. Weiglein, Comp. Phys. Comm. **124** (2000) 76; Also see <http://www.feynhiggs.de>.
9. M. Frank, S. Heinemeyer, W. Hollik and G. Weiglein, hep-ph/0212037. Also see [www.feynhiggs.de](http://www.feynhiggs.de).
10. S. Heinemeyer, W. Hollik and G. Weiglein, Eur. Phys. Jour. **C9** (1999) 343.
11. M. Carena, M. Quirós and C.E.M. Wagner, Nucl. Phys. **B461** (1996) 407.
12. M. Carena, S. Mrenna and C. Wagner, Phys. Rev. **D60** (1999) 075010.
13. ALEPH, DELPHI, L3 and OPAL Collaborations, OPAL TN699
14. A. Brignole, G. Degrassi, P. Slavich and F. Zwirner, Nucl. Phys. B **631** (2002) 195; G. Degrassi, S. Heinemeyer, W. Hollik, P. Slavich and G. Weiglein, Eur. Phys. J. C **28** (2003) 133.
15. D.E. Groom *et al.*, Eur. Phys. J. **C15** (2000) 1, available on the PDG WWW pages (URL: <http://pdg.lbl.gov/>).
16. M. Coca [CDF & D0 Collaborations], FERMILAB-CONF-03-238-E Presented at Flavor Physics and CP Violation (FPCP 2003), Paris, France, 3-6 Jun 2003
17. J. Fernandez [DELPHI Collaboration], DELPHI 2003-045-CONF-665, contributed paper no. 320
18. P. Achard *et al.* [L3 Collaboration], Phys. Lett. B **545** (2002) 30
19. M. Boonekamp, *Flavour and Model Independent Higgs Searches*, these proceedings
20. A. Heister *et al.* [ALEPH Collaboration], Phys. Lett. B **526** (2002) 191.